

**MODELLING NEGATIVE THERMOMECHANICAL EFFECTS  
IN REINFORCED ROAD STRUCTURES  
WITH THERMOELASTIC INCOMPATIBILITY OF COATING  
AND REINFORCEMENT MATERIALS**

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**Abstract.** The phenomena of the formation of local defects and cracks in asphalt concrete pavements of roads and bridges are most often observed in climatic zones with large temperature differences during their seasonal and daily changes. To a large extent, this is due to the heterogeneity of the thermomechanical properties of the materials of the coating layers and the base. To prevent these phenomena, reinforcing rods and meshes are introduced into the coating structure. In this work, using the theory of thermoelasticity, it is shown by the method of mathematical modelling that in cases of incompatibility of the thermomechanical characteristics of asphalt concrete materials and reinforcement, additional localized thermal stresses arise in its small vicinity, which, even at moderate temperatures, can reach critical values and lead to local defects and cracks. Since these defects are latent, they cannot always be detected in practice. The presented results of analytic calculation validated these conclusions. They can be used in both road building and composite design.

**Keywords:** reinforced asphalt concretes, thermomechanical incompatibility, mathematical modelling, destruction prevention.

**INTRODUCTION**

The strength and durability of the roadway is largely determined by the intensity of traffic loads and the impact of climatic conditions. Noticeable destruction of road surfaces, bridges, tunnels and dams, as well as other infrastructure facilities in climatic zones with large temperature differences, as a rule, occurs during off-season periods, accompanied by high temperature gradients.

Among the most common types of thermal destruction of the roadway is the appearance of transverse cracks in it, caused by the limiting values of longitudinal stresses at low negative temperatures in the conditions of the impossibility of free shortening of the upper layers. To avoid this effect, so-called “unloading expansion joints” and reinforcement (longitudinal, mesh, etc.) are introduced into the road structure. Such a general strengthening of the roadway with reinforcement leads to an increase in its overall strength, a reduction of deformability, an enlargement of durability, and a decrease in the cost of repair work.

In the theoretical analysis of the effect of reinforcement on the structural strength and the study of the general thermomechanical properties of reinforced (composite) materials and road coatings, the reduced (effective) values of the parameters of combined systems containing inclusions in the form of particles, fibers or rods are mainly determined [4, 6]. In these cases, mainly, models of homogeneous and inhomogeneous spherical particles, including those coated with shell layers, are considered [4]. The cases of ordered [2] and stochastic [10] placement of grains of these particles are singled out, and the reduced values of Young's modulus, Poisson's ratio, thermal conductivity coefficient, and thermal expansion coefficient of the entire system are calculated for them.

Very complex processes of thermal deformation and thermal destruction are observed in the structures of asphalt concrete pavements of roads and bridges [3, 9, 13, 16, 17, 19, 20]. The issues of determining the reduced thermomechanical characteristics of asphalt concrete materials reinforced with particles, fibers and rods are considered in publications [5, 7, 12, 14]. Here, however, these tasks become more complicated, since it is possible to create materials with directional (anisotropic) properties.

In addition, it should be noted that the insertion of reinforcing inclusions from another material into one material can not only improve the generalized characteristics of the entire composite, but under thermal effects it can also be accompanied by the generation of noticeable additional local internal thermal stresses if the thermomechanical characteristics of the composite components are incompatible. For plastic materials, these stresses can lead to local plastic deformations and defects; for brittle materials, to local cracking. Since these defects are localized and latent, they are not always detectable. Therefore, the problem of their theoretical forecasting seems to be relevant.

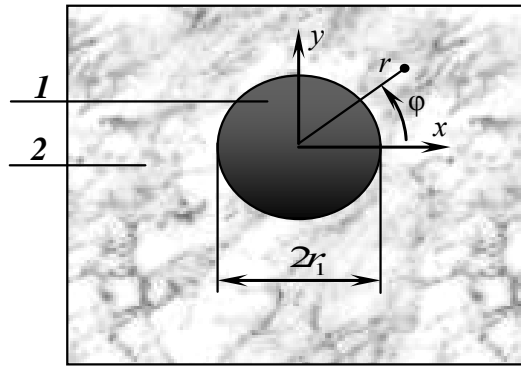
To simulate these effects, in this work, on the basis of the theory of thermoelasticity, the problem is posed of a planar thermally deformed state of an elastic medium containing an elastic rod of a circular cross section with different thermomechanical parameters. For the case of a change in the temperature of the system by a constant value, an analytical solution of the constitutive equations is constructed, expressions for thermal deformations and thermal stresses are obtained. The general regularities of possible negative influence of the thermomechanical incompatibility of the system parameters on the internal fields of the additional stresses are found. It has been established that the maximum thermal stresses in the medium are realized on the surface of its contact with an elastic inclusion, and they decrease along the radial coordinate in proportion to the square of the distance to the rod axis. The conditions for thermomechanical compatibility of the properties of the medium and the rod are formulated, under which there are no additional thermal stresses in the system. It is shown that in a system with incompatible parameters, additional thermal stresses can be decreased by reducing the radial rigidity of the inclusion through insertion a cylindrical cavity into it.

## **STATEMENT OF THE PROBLEM**

Let us formulate the problem of stationary thermal deformation of an infinite elastic medium 2 (matrix), which is reinforced with rod 1 of circular cross section of radius  $r_1$ . Fig. 1 shows a fragment of this system.

Let's use a cylindrical coordinate system  $O\varphi z$ , axis  $Oz$  of which coincides with the axis of the rod. Let the thermomechanical characteristics of rod 1 and medium 2 be determined, respectively, by the Lamé parameters  $\lambda_1, \mu_1$  and  $\lambda_2, \mu_2$  and coefficients of thermal linear expansion  $\alpha_1$  and  $\alpha_2$ . The temperature of the system changes steadily by the value  $\Delta T$ . Let us single out the case when the thermoelastic relative strains  $\varepsilon_z^{(i)}(r, \varphi, z)$  of the bodies  $i=1,2$  along the  $Oz$  axis are equal to zero and the system is in a plane axisymmetric thermally deformed state, described by the equilibrium equations [1, 8, 11, 15, 18]

$$\frac{d\sigma_r^{(i)}}{dr} + \frac{\sigma_r^{(i)} - \sigma_\varphi^{(i)}}{r} = 0, \quad (i=1,2), \quad (1)$$



*Fig. 1. Planar fragment of an elastic medium with a rod inclusion*

where  $\sigma_r^{(i)}, \sigma_\varphi^{(i)}$  are the normal radial and circumferential stresses of bodies 1 and 2 on the respective areas  $r = \text{const}$  and  $\varphi = \text{const}$ . Let us express normal thermal stresses in terms of strains  $\varepsilon_r^{(i)}, \varepsilon_\varphi^{(i)}, \varepsilon_z^{(i)}$ :

$$\begin{aligned} \sigma_r^{(i)}(r) &= (\lambda_i + 2\mu_i)\varepsilon_r^{(i)} + \lambda_i(\varepsilon_\varphi^{(i)} + \varepsilon_z^{(i)}) - (3\lambda_i + 2\mu_i)\lambda_i\Delta T; \\ \sigma_\varphi^{(i)}(r) &= (\lambda_i + 2\mu_i)\varepsilon_\varphi^{(i)} + \lambda_i(\varepsilon_r^{(i)} + \varepsilon_z^{(i)}) - (3\lambda_i + 2\mu_i)\alpha_i\Delta T; \\ \sigma_z^{(i)}(r) &= (\lambda_i + 2\mu_i)\varepsilon_z^{(i)} + \lambda_i(\varepsilon_r^{(i)} + \varepsilon_\varphi^{(i)}) - (3\lambda_i + 2\mu_i)\alpha_i\Delta T, \quad (i=1,2). \end{aligned} \quad (2)$$

Next, we take into account that  $\varepsilon_z^{(i)}(r, \varphi, z) = 0$ . Then expressions (2) will be simplified

$$\begin{aligned} \sigma_r^{(i)}(r) &= (\lambda_i + 2\mu_i)\varepsilon_r^{(i)} + \lambda_i\varepsilon_\varphi^{(i)} - (3\lambda_i + 2\mu_i)\alpha_i\Delta T; \\ \sigma_\varphi^{(i)}(r) &= (\lambda_i + 2\mu_i)\varepsilon_\varphi^{(i)} + \lambda_i\varepsilon_r^{(i)} - (3\lambda_i + 2\mu_i)\alpha_i\Delta T; \\ \sigma_z^{(i)}(r) &= \lambda_i(\varepsilon_r^{(i)} + \varepsilon_\varphi^{(i)}) - (3\lambda_i + 2\mu_i)\alpha_i\Delta T, \quad (i=1,2). \end{aligned} \quad (3)$$

The deformations used in (3) depend on the radial displacement  $u(r)$ :

$$\varepsilon_r^{(i)}(r) = \frac{\partial u^{(i)}}{\partial r}, \quad \varepsilon_\varphi^{(i)}(r) = \frac{u^{(i)}}{r}, \quad (i=1,2). \quad (4)$$

Taking into account (3), (4), equation (1) is reduced to the form

$$\frac{d^2 u^{(i)}}{dr^2} + \frac{1}{r} \frac{du^{(i)}}{dr} - \frac{1}{r^2} u^{(i)} = 0 \quad (5)$$

for each body  $i = 1, 2$ .

Let us represent equation (5) in a more compact form:

$$\frac{d}{dr} \left[ \frac{1}{r} \frac{d}{dr} (ru^{(i)}) \right] = 0, \quad (i = 1, 2). \quad (6)$$

Integrating the left side of equation (6) twice over  $r$ , get his solutions

$$u^{(1)}(r) = rC_1 + \frac{1}{r}C_2 \text{ at } i = 1;$$

$$u^{(2)}(r) = rC_3 + \frac{1}{r}C_4 \text{ at } i = 2. \quad (7)$$

The unknown constants  $C_i$  ( $i = \overline{1, 4}$ ) included here are determined from the boundary conditions and the contact equation for  $r = r_1$ :

$$u^{(0)}(0) = 0; \quad (8)$$

$$u_r^{(1)}(r_1) = u_r^{(2)}(r_1); \quad (9)$$

$$\sigma_r^{(1)}(r_1) = \sigma_r^{(2)}(r_1); \quad (10)$$

$$\sigma_r^{(2)}(r) \rightarrow 0 \text{ at } r \rightarrow \infty. \quad (11)$$

Condition (8) implies

$$C_2 = 0.$$

Using equalities (7), we express the strains and stresses of bodies 1 and 2 in terms of  $C_i$  ( $i = 1, 3, 4$ ):

$$\varepsilon_r^{(1)}(r) = C_1, \quad \varepsilon_\phi^{(1)}(r) = C_1; \quad \varepsilon_r^{(2)}(r) = C_3 - \frac{1}{r^2}C_4, \quad \varepsilon_\phi^{(2)}(r) = C_3 + \frac{1}{r^2}C_4;$$

$$\sigma_r^{(1)}(r) = 2(\lambda_1 + \mu_1)C_1 - (3\lambda_1 + 2\mu_1)\alpha_1\Delta T; \quad (12)$$

$$\sigma_r^{(2)}(r) = 2(\lambda_2 + \mu_2)C_3 - \frac{2\mu_2}{r^2}C_4 - (3\lambda_2 + 2\mu_2)\alpha_2\Delta T.$$

Condition (11) and the last equality of system (12) imply:

$$C_3 = \frac{3\lambda_2 + 2\mu_2}{2(\lambda_2 + \mu_2)}\alpha_2\Delta T.$$

Constants  $C_1$  and  $C_4$  are found from the system of equations (9), (10) transformed taking into account equalities (12),

$$C_1 = \frac{[(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 + (3\lambda_2 + 2\mu_2)\mu_2\alpha_2]\Delta T}{2(\lambda_1 + \mu_1 + \mu_2)(\lambda_2 + \mu_2)},$$

$$C_4 = \frac{r_1^2[(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 - (3\lambda_2 + 2\mu_2)(\lambda_1 + \mu_1)\alpha_2]\Delta T}{2(\lambda_1 + \mu_1 + \mu_2)(\lambda_2 + \mu_2)}.$$

Knowing constants  $C_i$  ( $i = \overline{1,4}$ ), find the displacement functions:

$$u^{(1)}(r) = \frac{r[(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 + (3\lambda_2 + 2\mu_2)\mu_2\alpha_2]\Delta T}{2(\lambda_1 + \mu_1 + \mu_2)(\lambda_2 + \mu_2)}, \quad (0 \leq r \leq r_1),$$

$$u^{(2)}(r) = r \frac{(3\lambda_1 + 2\mu_2)}{2(\lambda_2 + \mu_2)} \alpha_2 \Delta T +$$

$$+ \frac{r_1^2[(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 - (3\lambda_2 + 2\mu_2)(\lambda_1 + \mu_1)\alpha_2]\Delta T}{2(\lambda_1 + \mu_1 + \mu_2)(\lambda_2 + \mu_2)}, \quad (r \geq r_1).$$

Note that in the equation for  $u^{(2)}(r)$  the first term is the radial displacement in a homogeneous medium 2 in the absence of rod 1, the second term is due to the influence of body 1. It decreases in proportion to radius  $r$ .

We also give expressions for the stresses in rod 1:

$$\sigma_r^{(1)}(r) = \sigma_\phi^{(1)} = - \frac{\mu_2[(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 - (3\lambda_2 + 2\mu_2)(\lambda_1 + \mu_1)\alpha_2]\Delta T}{(\lambda_1 + \mu_1 + \mu_2)(\lambda_2 + \mu_2)};$$

$$\sigma_z^{(1)}(r) =$$

$$= \frac{[-(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)(\mu_1 + \mu_2)\alpha_1 + (3\lambda_2 + 2\mu_2)\lambda_1\mu_2\alpha_2]\Delta T}{(\lambda_1 + \mu_1 + \mu_2)(\lambda_2 + \mu_2)}, \quad (0 \leq r \leq r_1) \quad (13)$$

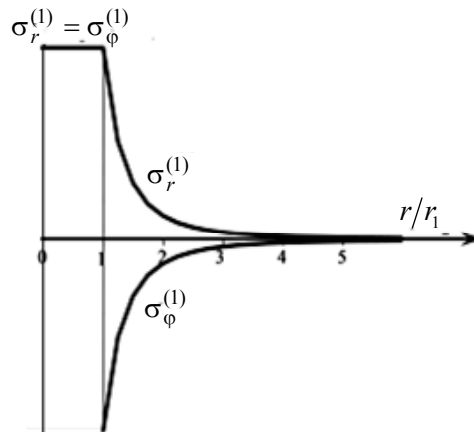
and in medium 2:

$$\sigma_r^{(2)}(r) = - \frac{r_1^2}{r^2} \frac{\mu_2[(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 - (3\lambda_2 + 2\mu_2)(\lambda_1 + \mu_1)\alpha_2]\Delta T}{(\lambda_1 + \mu_1 + \mu_2)(\lambda_2 + \mu_2)};$$

$$\sigma_\phi^{(2)}(r) = \frac{r_1^2}{r^2} \frac{\mu_2[(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 - (3\lambda_2 + 2\mu_2)(\lambda_1 + \mu_1)\alpha_2]\Delta T}{(\lambda_1 + \mu_1 + \mu_2)(\lambda_2 + \mu_2)};$$

$$\sigma_z^{(2)}(r) = - \frac{\mu_2(3\lambda_2 + 2\mu_2)}{(\lambda_2 + \mu_2)} \alpha_2 \Delta T, \quad (r \geq r_1). \quad (14)$$

Graphs of these functions are shown in Fig. 2. They indicate that additional



*Fig. 2. Graphs of the thermal stresses distribution in the plane of the axial section of the reinforced system under planar thermal deformation*

thermal stresses  $\sigma_r^{(2)}(r)$ ,  $\sigma_\phi^{(2)}(r)$ , caused by the inclusion of reinforcing rod 1 into medium 2, are local in nature and decrease in proportion to the square of the radial coordinate. In addition, they are equal to each other in absolute value and differ in signs, which depend on the ratio of quantities  $\alpha_1$  and  $\alpha_2$ . So if  $\alpha_1 > \alpha_2$ , then, as follows from the form of the numerators of formula (14) with  $\Delta T > 0$ ,  $\alpha_1 > \alpha_2$ , there are inequalities  $\sigma_r^{(2)}(r) < 0$ ,  $\sigma_\phi^{(2)}(r) > 0$ , and if  $\alpha_1 < \alpha_2$ , then vice versa,  $\sigma_r^{(2)}(r) > 0$ ,  $\sigma_\phi^{(2)}(r) < 0$ . This means that, since asphalt concrete has a lower tensile strength than compressive strength, under any temperature  $\Delta T$  changes, unfavorable thermal stresses will be realized for  $\sigma_r^{(2)}(r)$  or  $\sigma_\phi^{(2)}(r)$ .

It is also obvious that the values of thermomechanical parameters, at which the numerators of fractions (14) of functions  $\sigma_r^{(2)}(r)$ ,  $\sigma_\phi^{(2)}(r)$  are zero, are thermally compatible. Therefore, equality

$$(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 - (3\lambda_2 + 2\mu_2)(\lambda_1 + \mu_1)\alpha_2 = 0 \quad (15)$$

represents a condition for the compatibility of the thermomechanical parameters of the matrix and the reinforcing rod.

Condition (15) can be simplified if to replace the Lamé parameters  $\lambda$  and  $\mu$  with modulus of elasticity  $E$  and Poisson's ratio  $\nu$ , using formulas

$$\lambda = \frac{E\nu}{(1+\nu)(1-2\nu)}; \quad \mu = \frac{E}{2(1+\nu)}.$$

Then, instead of (15) we have a simpler record of this condition

$$\frac{\alpha_1}{1+\nu_2} = \frac{\alpha_2}{1+\nu_1}.$$

As an example, consider the case when a fiberglass reinforcing rod of radius  $r_1$  with thermomechanical parameters  $\lambda_1 = 30.2$  GPa,  $\mu_1 = 12.9$  GPa,  $\alpha_1 = 21 \cdot 10^{-6} \text{ K}^{-1}$  is located in an asphalt concrete medium with parameters  $\lambda_2 = 1.39$  GPa,  $\mu_2 = 2.08$  GPa,  $\alpha_2 = 10 \cdot 10^{-6} \text{ K}^{-1}$ . It is accepted that  $\Delta T = -20^\circ \text{ K}$ . With these data, the thermal stresses in the system amounted to  $\sigma_r^{(1)} = \sigma_\phi^{(1)} = 1.298$  MPa,  $\sigma_z^{(1)} = 15.541$  MPa,  $\sigma_r^{(2)} = 1.298 \frac{r_1^2}{r^2}$  MPa,  $\sigma_z^{(2)} = -1.298 \frac{r_1^2}{r^2}$  MPa,  $\sigma_\phi^{(2)} = 0.9986$  MPa.

If we take into account that the ultimate strength of asphalt concrete in compression is  $5 \div 20$  MPa, and in tension it turns out to be several times less than these values, then we can conclude that under the considered conditions, additional thermal stresses in asphalt concrete, caused by the insertion of a fiberglass reinforcing rod into it, can lead to the occurrence of local defects in its small neighbourhood.

### REDUCING THE LEVEL OF ADDITIONAL THERMAL STRESSES ON TUBULAR REINFORCING RODS

As can be seen from equalities (13), (14), additional stresses  $\sigma_r^{(i)}$ ,  $\sigma_\phi^{(i)}$  ( $i=1,2$ ) are determined not only by the values of coefficients  $\alpha_1$ ,  $\alpha_2$ , but also by elasticity parameters  $\lambda_i$ ,  $\mu_i$  ( $i=1,2$ ), which are included in the numerators of fractions (13), (14) in the third powers, and in the denominators — in the second ones. Therefore, additional thermal stresses in the system increase with increasing  $\lambda_i$ ,  $\mu_i$  ( $i=1,2$ ) or, for example, with an increase in the radial rigidity of rod 1 (while maintaining its axial strength and rigidity). Conversely, they decrease as this stiffness decreases. Given this property, we can propose to use tubular rods as reinforcement in asphalt concrete pavements (Fig. 3). Let us investigate the thermally stressed state in this case. Let  $r_1$  and  $r_2$  be the inner and outer radii of pipe 1, respectively,

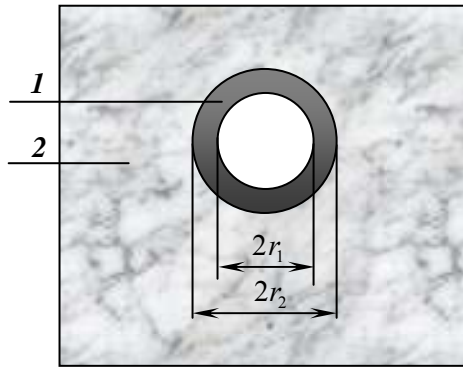


Fig. 3. Planar fragment of an elastic medium with a tubular inclusion

the dimensions of medium 2 are unlimited. Let us assume, as above, that  $\lambda_i$ ,  $\mu_i$ ,  $\alpha_i$  ( $i=1,2$ ) be the thermomechanical parameters of bodies 1 and 2,  $\Delta T$  — difference in body temperature in the initial and final states. Let us find the functions of thermal stresses in the system for the case of its axisymmetric planar thermally deformed state.

Similarly to the case of a solid rod, the equations of thermoelasticity of the system have form (1) – (6). The solution of these equations is again formulated in

the form of functions of radial displacements of body 1

$$u_1(r) = rC_1 + \frac{1}{r}C_2 \quad (r_1 \leq r \leq r_2),$$

and for medium 2

$$u_2(r) = rC_3 + \frac{1}{r}C_4 \quad (r \geq r_2).$$

At the same time, constants  $C_i$  ( $i=1,4$ ) are found from the conditions:

$$\sigma_r^{(1)}(r_1) = 0; \quad u^{(1)}(r_2) = u^{(2)}(r_2);$$

$$\sigma_r^{(1)}(r_2) = \sigma_r^{(2)}(r_2); \quad \sigma_r^{(2)}(r) \rightarrow 0 \quad \text{at } r \rightarrow \infty.$$

After appropriate substitutions, these equations are reduced to the form:

$$2(\lambda_1 + \mu_1)C_1 - \frac{2\mu_1}{r_1^2}C_2 - (3\lambda_1 + 2\mu_1)\alpha_1\Delta T = 0;$$

$$C_1 + \frac{1}{r_2^2}C_2 - C_3 - \frac{1}{r_2^2}C_4 = 0;$$

$$\begin{aligned}
 & 2(\lambda_1 + \mu_1)C_1 - \frac{2\mu_1}{r_2^2}C_2 - (3\lambda_1 + 2\mu_1)\alpha_1\Delta T - \\
 & - 2(\lambda_2 + \mu_2)C_3 + \frac{2\mu_2}{r_2^2}C_4 + (3\lambda_2 + 2\mu_2)\alpha_2\Delta T = 0; \\
 & 2(\lambda_2 + \mu_2)C_3 - (3\lambda_2 + 2\mu_2)\alpha_2\Delta T = 0.
 \end{aligned} \tag{16}$$

From the last equation of this system, we obtain

$$C_3 = \frac{(3\lambda_2 + 2\mu_2)}{2(\lambda_2 + \mu_2)}\alpha_2\Delta T.$$

Next, from the remaining equations of system (16) we find:

$$\begin{aligned}
 C_1 &= \frac{\mu_1}{r_1^2} \frac{[-(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 + (3\lambda_2 + 2\mu_2)(\lambda_1 + \mu_1)\alpha_2]\Delta T}{2(\lambda_1 + \mu_1)^2(\lambda_2 + \mu_2) \left[ \frac{\mu_1}{r_1^2(\lambda_1 + \mu_1)} + \frac{1}{r_2^2} + \frac{\mu_1}{\mu_2} \left( \frac{1}{r_1^2} - \frac{1}{r_2^2} \right) \right]} + \\
 & + \frac{(3\lambda_1 + 2\mu_1)}{2(\lambda_1 + \mu_1)}\alpha_1\Delta T; \\
 C_2 &= \frac{[-(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 + (3\lambda_2 + 2\mu_2)(\lambda_1 + \mu_1)\alpha_2]\Delta T}{2(\lambda_1 + \mu_1)(\lambda_2 + \mu_2) \left[ \frac{\mu_1}{r_1^2(\lambda_1 + \mu_1)} + \frac{1}{r_2^2} + \frac{\mu_1}{\mu_2} \left( \frac{1}{r_1^2} - \frac{1}{r_2^2} \right) \right]}, \\
 C_4 &= \frac{r_2^2\mu_1}{\mu_2} \left( \frac{1}{r_1^2} - \frac{1}{r_2^2} \right) \frac{[(3\lambda_1 + 2\mu_1)(\lambda_2 + \mu_2)\alpha_1 - (3\lambda_2 + 2\mu_2)(\lambda_1 + \mu_1)\alpha_2]\Delta T}{2(\lambda_1 + \mu_1)(\lambda_2 + \mu_2) \left[ \frac{\mu_1}{r_1^2(\lambda_1 + \mu_1)} + \frac{1}{r_2^2} + \frac{\mu_1}{\mu_2} \left( \frac{1}{r_1^2} - \frac{1}{r_2^2} \right) \right]}.
 \end{aligned}$$

Using the found constants, you can build expressions for displacements

$$u(r) = rC_1 + \frac{1}{r}C_2, \quad (r_1 \leq r \leq r_2); \quad u(r) = rC_3 + \frac{1}{r}C_4, \quad (r \geq r_2)$$

and stresses

$$\sigma_r^{(1)}(r) = 2(\lambda_1 + \mu_1)C_1 - \frac{2\mu_1}{r^2}C_2 - (3\lambda_1 + 2\mu_1)\alpha_1\Delta T;$$

$$\sigma_\phi^{(1)}(r) = 2(\lambda_1 + \mu_1)C_1 + \frac{2\mu_1}{r^2}C_2 - (3\lambda_1 + 2\mu_1)\alpha_1\Delta T;$$

$$\sigma_z^{(1)}(r) = 2\lambda_1C_1 - (3\lambda_1 + 2\mu_1)\alpha_1\Delta T, \quad (r_1 \leq r \leq r_2)$$

in rod 1 and

$$\sigma_r^{(2)}(r) = -\frac{2\mu_2}{r^2}C_4; \quad \sigma_\phi^{(2)}(r) = -\frac{2\mu_1}{r^2}C_2;$$

$$\sigma_z^{(2)}(r) = 2\lambda_2C_3 - (3\lambda_2 + 2\mu_2)\alpha_2\Delta T, \quad (r \geq r_1)$$

in medium 2.

Table shows the stress values  $\sigma_r^{(1)}$ ,  $\sigma_\phi^{(1)}$ ,  $\sigma_z^{(1)}$  on surfaces  $r = r_1$  and  $r = r_2$  of fiberglass body 1 and stresses  $\sigma_r^{(2)}$ ,  $\sigma_\phi^{(2)}$ ,  $\sigma_z^{(2)}$  on surface  $r = r_2$  of asphalt



concrete medium 2 for relations  $r_1/r_2 = 0.5, 0.75,$  and  $0.9$  at the values of the thermomechanical parameters given above and the temperature difference  $\Delta T = -20$  K.

Values of thermal stresses in the medium reinforced with a tubular rod

Types of thermal stresses	$r_1/r_2$			
	0.0	0.5	0.75	0.9
$\sigma_r^{(1)}(r_1)$ MPa	1.2981	0	0	0
$\sigma_\phi^{(1)}(r_1)$ MPa	1.2981	3.2455	4.5766	7.3764
$\sigma_z^{(1)}(r_1)$ MPa	15.541	15.569	16.235	17.215
$\sigma_r^{(1)}(r_2)$ MPa	1.2981	1.2171	1.0011	0.7008
$\sigma_\phi^{(1)}(r_2)$ MPa	1.2981	2.0285	3.5742	6.6752
$\sigma_z^{(1)}(r_2)$ MPa	15.541	15.569	16.235	17.215
$\sigma_r^{(2)}(r_2)$ MPa	1.2981	1.2171	1.0011	0.7008
$\sigma_\phi^{(2)}(r_2)$ MPa	-1.2981	-1.2171	-1.0011	-0.7008
$\sigma_z^{(2)}(r_2)$ MPa	0.9986	0.9986	0.9986	0.9986

The question of the distribution in the radial direction of additional thermal stresses  $\sigma_r^{(i)}(r), \sigma_\phi^{(i)}(r)$  deserves a special interest. Fig. 4 shows the graphs of these functions for case  $r_1/r_2 = 0.75$ .

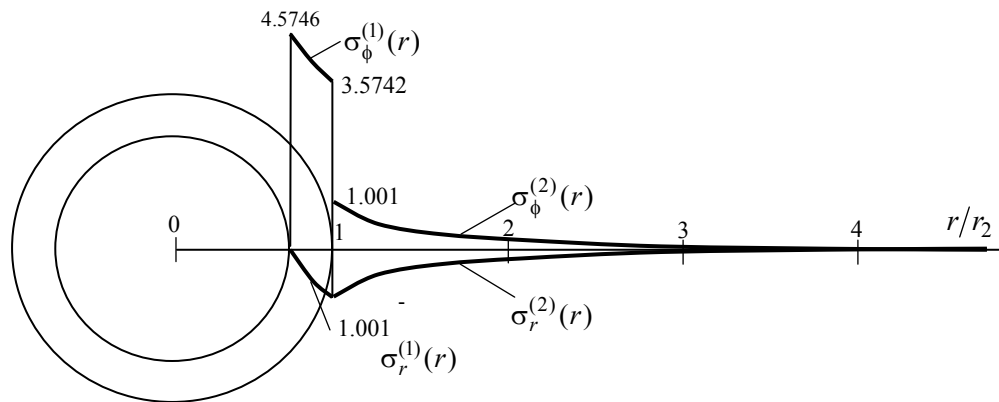


Fig. 4. Graphs of distribution of thermal stresses  $\sigma_r^{(i)}(r), \sigma_\phi^{(i)}(r)$  (MPa) for case  $r_1/r_2 = 0,75$

As can be seen, additional thermal stresses  $\sigma_r^{(2)}(r), \sigma_\phi^{(2)}(r)$  in the asphalt concrete have the highest values on the contact surface  $r = r_2$  and they decrease rapidly in the radial direction.

An analysis of the above results indicates that the replacement of a solid reinforcing rod with a tubular one leads to a noticeable decrease in additional local thermal stresses in the matrix medium, although the thermal stresses in the rod increase somewhat. This effect becomes more noticeable as the thickness of the tube rod decreases.

## CONCLUSION

1. The problem associated with the modelling of the formation of additional thermal stresses, defects and destructions in the medium of an asphalt concrete pavement with the insertion of a reinforcing rod into it is considered. On the basis of thermoelasticity methods, the system of differential equations is formed for a plane axisymmetric deformed state of an infinite cylindrical elastic body in an infinite elastic medium under condition of a change in the temperature of the system.

2. An analytical solution of the formulated equations is constructed in a closed form, which determines the additional thermal displacements, additional strains and stresses in the system. It is shown that additional thermal stresses have the highest values on the contact surface of the reinforcing rod and the coating array and decrease in inverse proportion to the square of the distance from this surface. It has been established that the values of these stresses enlarge with an increase in the thermomechanical incompatibility of the system materials and the radial stiffness of the rod. On the example of asphalt concrete reinforced with a fiberglass rod, it was demonstrated that even with moderate temperature changes, additional thermal stresses in asphalt concrete can reach critical values.

3. A method is proposed for reducing additional thermal stresses by reducing the radial stiffness of the reinforcement by replacing a solid rod with a hollow tube. Theoretical modelling of this effect showed that with a decrease in the tube wall thickness, the decrease in additional contact thermal stresses in asphalt concrete becomes significant.

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#### НЕГАТИВНІ ТЕРМОМЕХАНІЧНІ ЕФЕКТИ В АРМОВАНИХ ДОРОЖНІХ КОНСТРУКЦІЯХ ЗА ТЕРМОПРУЖНОЇ НЕСУМІСНОСТІ МАТЕРІАЛІВ ПОКРИТТЯ ТА АРМАТУРИ / В.І. Гуляєв, В.В. Мозговий, Н.В. Шлюнь, Л.В. Шевчук

**Анотація.** Явища утворення локальних дефектів і тріщин в асфальтобетонних покриттях автомобільних доріг та мостів найчастіше спостерігаються у кліматичних зонах з великими перепадами температур за їх сезонних та добових змін. Значною мірою це зумовлено неоднорідністю термомеханічних властивостей матеріалів шарів покриттів та основи. Для попередження цих явищ в конструкції покриттів вводять армувальні стрижні і сітки. У роботі методами теорії термопружності показано, що у випадках несумісності термомеханічних характеристик матеріалів асфальтобетону та арматури в її малому околі виникають додаткові локалізовані термонапруження, які навіть за помірних значень перепадів температури можуть досягати критичних значень та призводити до локальних дефектів і тріщин. Оскільки ці дефекти мають прихований характер, їх не завжди можна виявляти.

**Ключові слова:** асфальтобетонне покриття, стрижнева арматура, термомеханічна несумісність, концентрація термонапруг.